

Die Untersuchungen sollen an anderen N-Aryl-carbazolen fortgeführt werden.

8 Das gilt für Wellenlängen > 230 mµ. Mit kürzeren Wellenlängen wurde nicht angeregt.

Scattering Cross Section of a Particle Beam

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Using usual methods the scattering cross section of a beam of wave packets is derived.

The scattering cross section for a beam of wave packets was shown to be the superposition of sections of plane wave components of the wave packet. We derive this interesting result in a more usual way. The transition probability per unit time into the state defined by quantum numbers \( (a, E_a) \) is (see e. g. Newton)

\[
\begin{align*}
  w_a(t) &= -i \int \int dE_b \int dE_c f(E_b, b) f^*(E_c, c) \tau_{ab}^{(c)} \tau_{ab}^{(c)*} \exp(i(E_c - E_b) t) \\
  & \times \exp \left( \frac{(E_c - E_a - i \epsilon)^{-1} - (E_b - E_a + i \epsilon)^{-1}}{\epsilon} \right)
\end{align*}
\]

(1)

Here \( \tau_{ab}^{(c)} \) is the off-the energy shell extension of the transition matrix,

\[
\psi_a(t) = \int \int dE_b \exp(-iE_b t) \psi_b(E_b, b)
\]

is the incoming wave packet.

The total transition probability is \( W_a'' = \int dt w_a(t) \),

\[
W_a'' = (2 \pi)^2 \int dE_b f(E_b, b) f^*(E_a, c) \tau_{ab} \tau_{ac} \tau_{ac}^*(E_a)
\]

(2)

An "usual" packet is described by \( f(E_b, b) = mk_b f(k_b) \).

The transition probability into the state with the moment in an element of solid angle \( d\Omega \) is

\[
W_a' = \int dE_a \varrho(k_a) W_a''
\]

where with the wave function

\[
\psi_b(x | E_a, a) = \exp(i k_a x)
\]

the density of states is

\[
\varrho(k_a) = \frac{1}{(2 \pi)^3} k_a^2 d\Omega \Delta E_a.
\]

The z-axis is chosen along some median flux direction. The beam is considered to be composed of packets displaced relative to each other in the xy-plane. Denoting the displacement by \( h \), we obtain

\[
\tau^{(c)}_{h,k_b}(h) = \tau^{(c)}_{k_b,k_c} \exp(i k_c h).
\]

Then the distribution of the total flux in the z-axis direction is homogeneous

\[
j_z = \int dE_a \int d\Omega \int dt j_z(h) = (2 \pi)^2 \frac{d\Omega}{g(k)} |f(k)|^2.
\]

(3)

Similar calculations lead to the expression for the transition probability

\[
W_a'' = \int dE_a \int d\Omega \int dE_b |f(k)|^2 \tau_{h,k_b}(h)
\]

(4)

In both (3) and (4) eqs. the back scattered wave was neglected. The cross section is defined as
\[ \sigma_a = W_a j_2^{-1}. \] (5)
The cross section of the \( k \)-component of the wave packet is
\[ \sigma_a(k) = (m/2 \pi)^2 |\tau_{k,k}(k)|^2 d\Omega_a. \] (6)

Substituting (3), (4) and (6) into (5) and supposing \( \int dk |j(k)|^2 = 1 \), we obtain
\[ \sigma_a = \int \frac{dk}{k_2} |j(k)|^2 \sigma_a(k). \] (7)

Our expression is identical with Katz's (3.7)\(^1\): the scattering cross section of a beam of wave packets produced by an incoherent extended source is given by the superposition of the cross sections of \( k \)-components of the wave packet.

Chemical Reactions Resulting from Heavy Particle Bombardment of Diamond

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Irradiation of natural diamond crystals with protons and deuterons ranging in energy from 0.7 to 1.5 MeV has been found to produce infrared absorption bands which are apparently related to the formation of hydrocarbons. The process is essentially analogous to the formation of hydroxyl ions by proton irradiation of silicates\(^1\).

In the experiments reported here, the particle beam was supplied by a model Km, 3 MeV Van de Graaff accelerator manufactured by High-Voltage Engineering Inc. The irradiations varied in duration from two to twenty-seven hours with an average current of 1.5 \( \mu A \) on the target holder. Approximately \( 5 \times 10^{12} \) particles per hour were incident upon the crystal. In all cases the crystal was smaller than the beam diameter, and all irradiations were made with a de-focused beam so that the proton intensity distribution across the beam was as uniform as possible. The proton energy range between 0.7 and 1.5 MeV was selected in order to reduce nuclear reactions and to insure a depth of penetration which would permit retention of the reaction products.

Natural diamonds were found to be the most satisfactory target material because they are transparent to infrared radiation throughout most of the spectrum where hydrocarbon absorption bands would be expected. The diamonds used for the study were flat cleavage plates of approximately 1 mm thickness, having surface areas of 10 to 15 mm\(^2\). They were mounted in lead plates or imbedded in Wood's metal in copper holders and were in thermal contact with a dewar containing liquid nitrogen. The targets were cooled in order to retard escape of any of the volatile hydrocarbons which might be produced.

Very strong blue luminescence was noted at the beginning of each of the initial irradiations. This luminescence decreased rapidly in intensity but did not disappear entirely for approximately an hour. In several instances the diamonds were found to be in poor thermal contact with the liquid nitrogen dewar. In these cases, the luminescence changed very rapidly from a strong blue to a dark green. The diamonds were colored black in visible light in all cases, irrespective of the nature of the bombarding particles.

The infrared absorption measurements were made using a Perkin-Elmer model 421 infrared spectrophotometer. Absorption curves of the diamond crystals were measured before irradiation and after each subsequent particle dose. After each irradiation and before measurement of the infrared absorption curve, the surfaces of the diamonds were cleaned with steel and glass wool. Crystals were then examined under a microscope to make sure that any external deposits which had been transported to the target by the ion beam had been removed. Differential curves were obtained by subtracting the normalized infrared absorption curve of the irradiated crystal from the normalized absorption curve of the same crystal before irradiation.

An initial irradiation resulted in the development of a series of absorption lines which increased in intensity and definition when the total proton dose was increased by successive irradiations. Differential absorption spectra of proton and deuteron irradiated diamonds are shown in Fig. 1. The absorption peaks which are shown in the proton and deuteron differential curves include fundamental absorption frequencies and their overtones and combinations which are characteristic of CH and CD bonds\(^2\). The deuteron irradiations were made in order to test the hypothesis that the infrared absorption lines which appeared after proton irradiations resulted from hydrocarbons formed in the crystal. Owing to the greater mass of the deuterium nucleus, the absorptions which result from deuterated hydrocarbons are shifted in wavelength.

A further test using \(^3\)He nuclei was conducted to eliminate the possibility that the absorption peaks resulted from radiation induced defects and not from hydrocarbons. If radiation induced defects caused the absorptions, similar peaks should be developed irrespective of the nature of the bombarding particles. The \(^3\)He nuclei blackened the crystals, as in the case of the other radiations; but absorption bands were not developed at wavelengths of the principal CH or CD absorptions.
